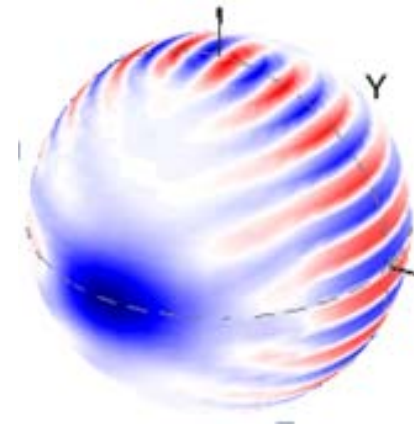
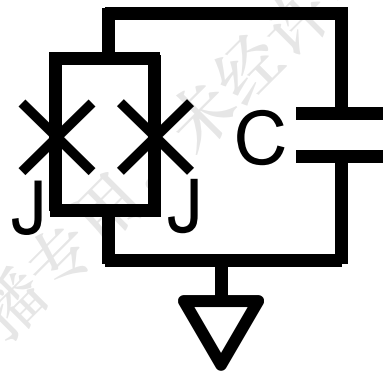
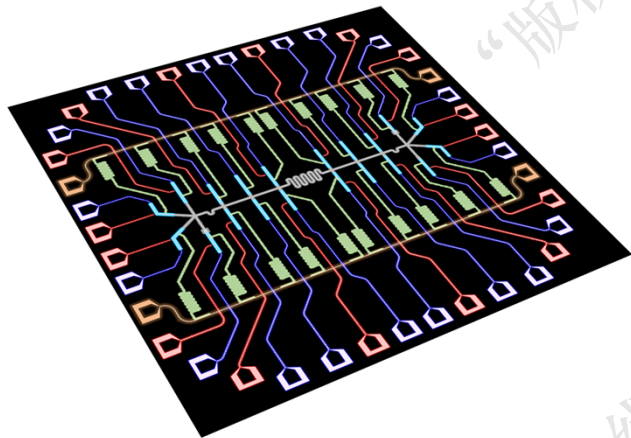


Superconducting circuits: A new platform for quantum optics

Shiyao Zhu, Da-Wei Wang & Haohua Wang

Department of Physics
Zhejiang University, Hangzhou, China



5th West-Lake Photonics Symposium, Nov. 28, 2020, Zhejiang

Outline

(I) Introduction

Fundamentals of superconducting qubits

(II) Experimental platform at Zhejiang University

(III) Quantum Optics Experiments with Superconducting-Qubits

1. Quantum chiral rotation

2. 5-components Schrodinger cat states

3. Control of superradiant and subradiant states

4. Excitation of two atoms with entangled photons

Important Development Related to Quantum information

1984: BB84 Scheme, Bennett and Brassard,

1993: Teleportation, Bennett, PRL 70, 1895

1997: teleportation experiment, A. Zeilinger, Nature 390, 575

1998: Kimble, Science 282,706

1998: D. Boschi et al, PRL 80, 1121

1985: Josephson joints, Clarke et al, PRL 55, 1543

1994: Shor Algorithm, P.W. Shor,

2001 IBM did the experiment factored 15 into 3×5 , Nature

1996: Grover Search, L.K. Grover,

1999: Made Josephson joints, Tsai et al, Nature 398, 786

2004: Josephson joints +Cavity, Schoelkopf et al, PRA 69, 062320

2007: Josephson joints with very long coherent time

J. Q. You (游建强), PRB 2007/4; and J. Koch, PRA 2007/10)

1945 (1930): Atomic clock (2000年1秒), I. Rabi (1944 Nobel Prize)

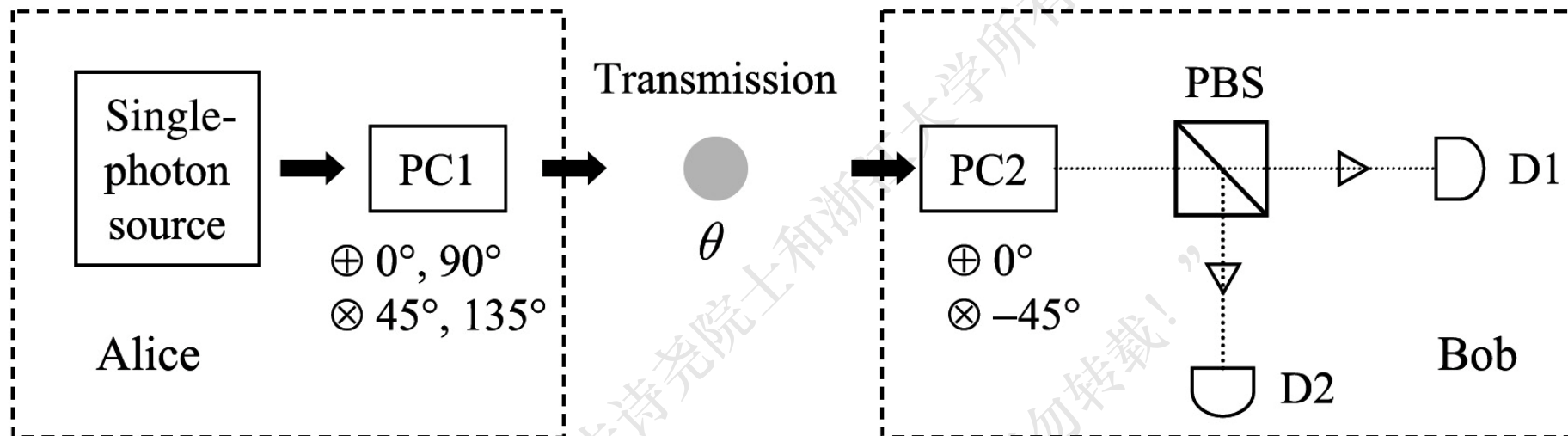
1955: First atomic clock, Louis Essen and Jack Parry

1989年Nobel Prize, Ramsey

Optic Clock



BB84 Schem: quantum key distribution



Alice's bit	0	1	1	0	1	0	0	1
Alice's basis	+	+	X	+	X	X	X	+
Alice's polarization	↑	→	↖	↑	↖	↗	↗	→
Bob's basis	+	X	X	X	+	X	+	+
Bob's measurement	↑	↗	↖	↗	→	↗	→	→
Public discussion								
Shared Secret key	0		1			0		1

Bennett and Brassard suggested in 1984年

Proc. Int. Conf. Comput. Syst. Signal Process



Shor算法：实验演示

NATURE | VOL 414 | 20/27 DECEMBER 2001

Experimental realization of Shor's quantum factoring algorithm using nuclear magnetic resonance

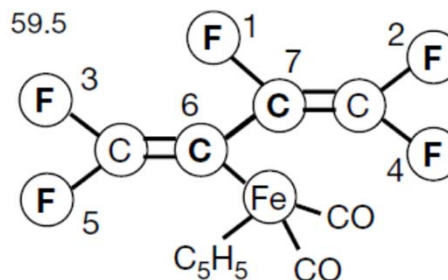
Lieven M. K. Vandersypen^{*†}, Matthias Steffen^{*†}, Gregory Breyta^{*}, Costantino S. Yannoni^{*}, Mark H. Sherwood^{*} & Isaac L. Chuang^{*†}

^{*} IBM Almaden Research Center, San Jose, California 95120, USA

[†] Solid State and Photonics Laboratory, Stanford University, Stanford, California 94305-4075, USA

7量子比特系统：
 $15 = 3 \times 5$

i	$\omega_i/2$	$T_{1,i}$	$T_{2,i}$	J_{7i}	J_{6i}	J_{5i}	J_{4i}	J_{3i}	J_{2i}
1	-22052.0	5.0	1.3	-221.0	37.7	6.6	-114.3	14.5	25.16
2	489.5	13.7	1.8	18.6	-3.9	2.5	79.9	3.9	
3	25088.3	3.0	2.5	1.0	-13.5	41.6	12.9		
4	-4918.7	10.0	1.7	54.1	-5.7	2.1			
5	15186.6	2.8	1.8	19.4	59.5				
6	-4519.1	45.4	2.0	68.9					
7	4244.3	31.6	2.0						





Grover's algorithm: quantum search

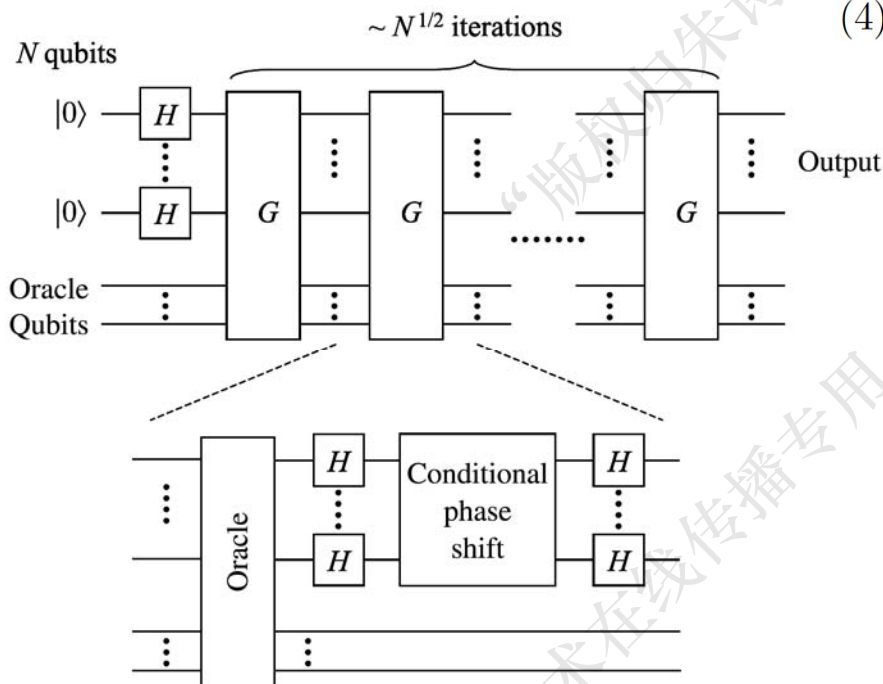
L.K. Grover: *A fast quantum mechanical algorithm for database search*,
 Proceedings, 28th Annual ACM Symposium on the Theory of Computing, p. 212, May **1996**.

$$|\psi_1\rangle = \left(\frac{1}{\sqrt{2}}\right)^N (|0\rangle_1 + |1\rangle_1)(|0\rangle_2 + |1\rangle_2) \cdots (|0\rangle_N + |1\rangle_N)$$

$$= \frac{1}{N_{\text{data}}^{1/2}} \sum_{x=0}^{N_{\text{data}}-1} |x\rangle,$$

The Grover operator employs four steps:

- (1) apply the **oracle** operator;
- (2) apply a Hadamard gate to each qubit in the register;
- (3) apply a conditional phase shift;
- (4) apply a Hadamard gate to each qubit in the register.



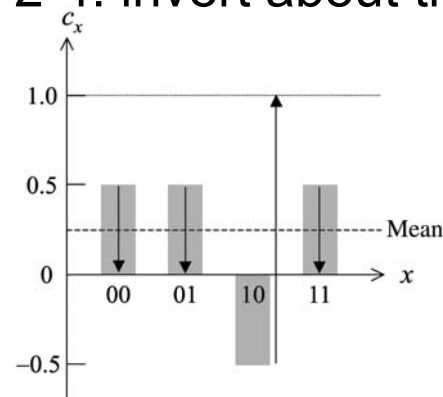
oracle:

$$|x\rangle \rightarrow (-1)^{f(x)}|x\rangle,$$

where $f(x)$ is a function defined by:

$$\begin{aligned} f(x) &= 1 && \text{if } x \text{ is the solution,} \\ f(x) &= 0 && \text{otherwise.} \end{aligned}$$

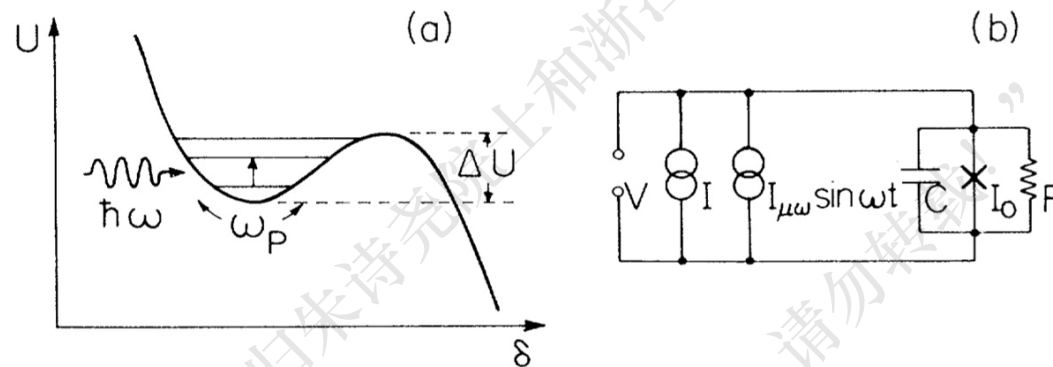
steps 2-4: invert about the mean



Superconducting quantum qubits

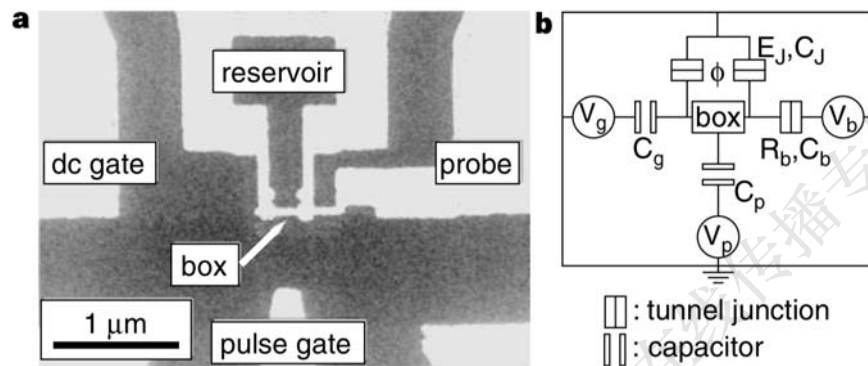
Josephson joints + Cavity

Quantized levels of Josephson junction



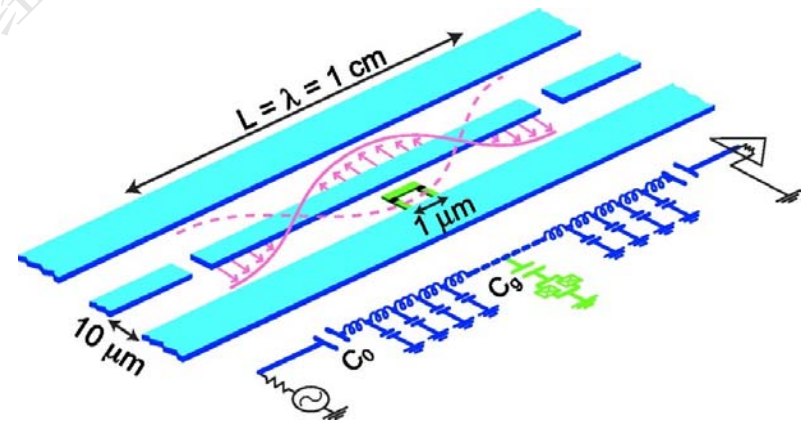
Clarke et al, PRL 55, 1543 (1985)

Qubit in a Cooper pair box



Tsai et al, Nature 398, 786 (1999)

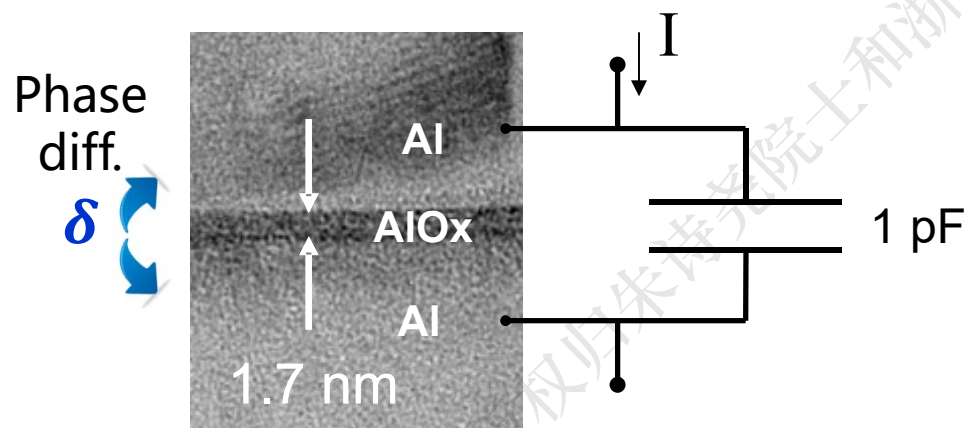
Cavity QED with coplanar resonator



Schoelkopf et al, PRA 69, 062320 (2004)

Josephson junction—the heart of a qubit

- Josephson junction:
nonlinear inductor

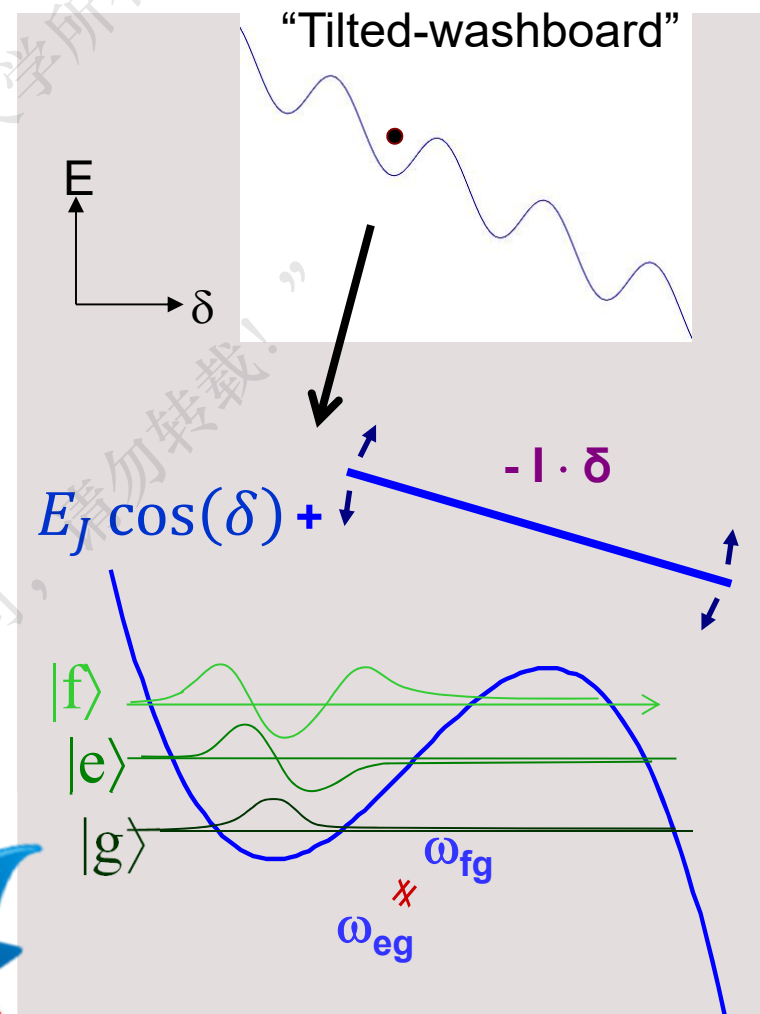


$$I_J = I_c \sin(\delta), \quad V = \frac{\Phi_0}{2\pi} \frac{d\delta}{dt}$$

Nonlinear $\hat{L}C$ resonator:

Stores quantized EM fields

$|g\rangle$ (low energy), $|e\rangle \rightarrow |0\rangle, |1\rangle$ of a qubit

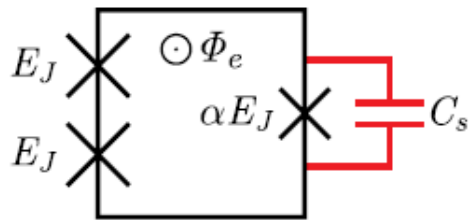


Artificial atom:
level separation ~ 6 GHz,
 ~ 2 GHz tunability

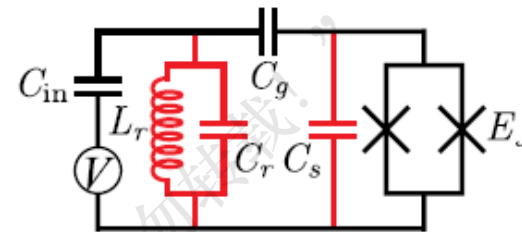
Qubit with reduced noise

Coherent time: 0.01ms - 0.1ms; 3-4 order longer.

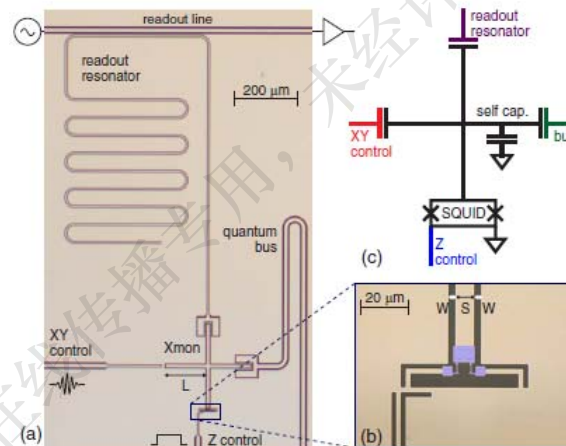
C-shunt flux qubit
J. Q. You et al., PRB **2007**.



Transmon qubits (C-shunt Cooper-pair box)
J. Koch et al., PRA **2007** (Yale组).



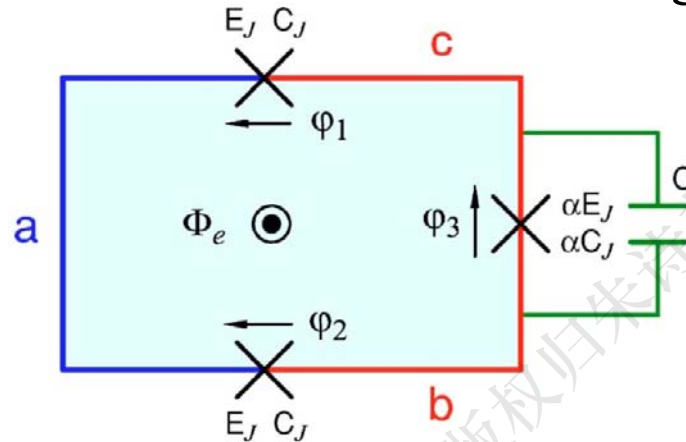
Xmon qubits: similar to deformed Transmon
R. Barends et al., PRL 2013 (UCSB/Google组).



Make quantum supremacy and quantum simulation possible.

Qubit with reduced noise

Make the coherent time 3 order longer

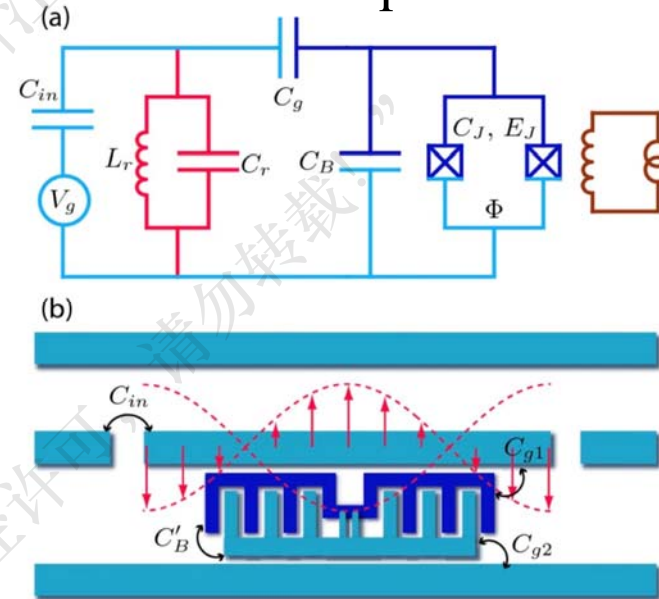


C-shunt flux qubit

You et al, PRB 75, 140515 (2007)

(Received 3 April 2007; published 27 April 2007)

Transmon qubits

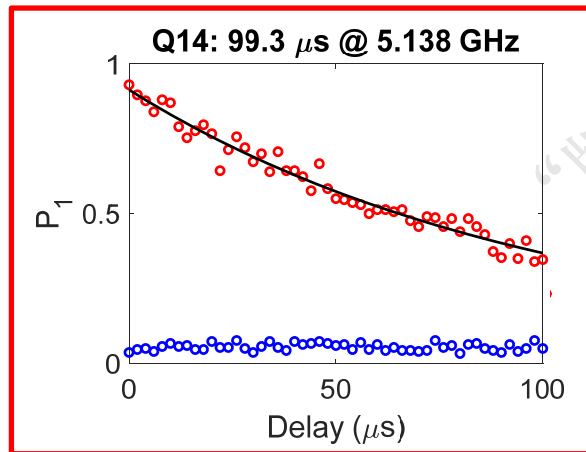
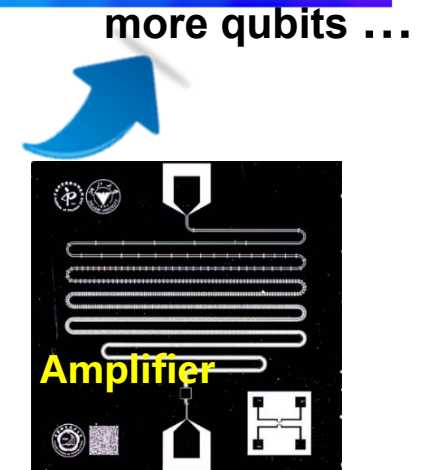
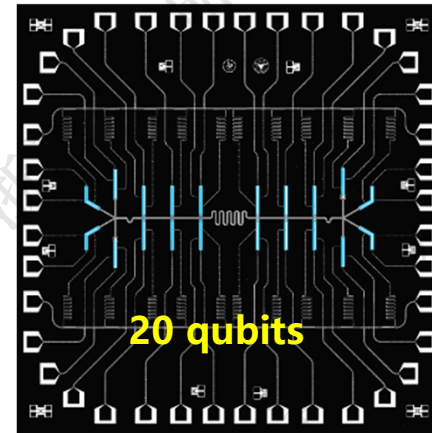
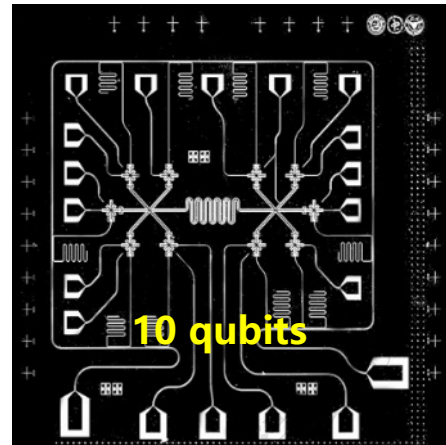
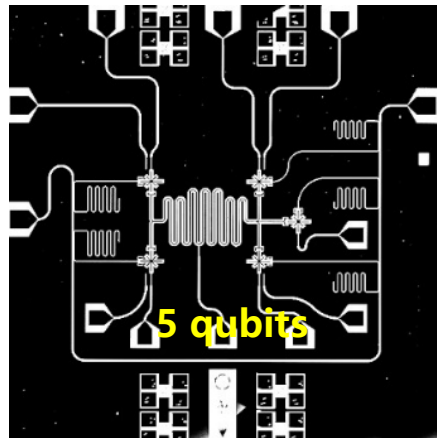


Koch et al, PRA 76, 042319 (2007)

(Received 22 May 2007; published 12 October 2007)

(II) Superconducting quantum qubits as an experiment platform in Zhejiang University

Characterizations



- ✓ **Freq. tunable** qubits: Best T_1 close to 100 μs @ 5GHz,
about 20 – 50 μs on average,
 T_2^* above 10 μs (5 – 10 μs on average).

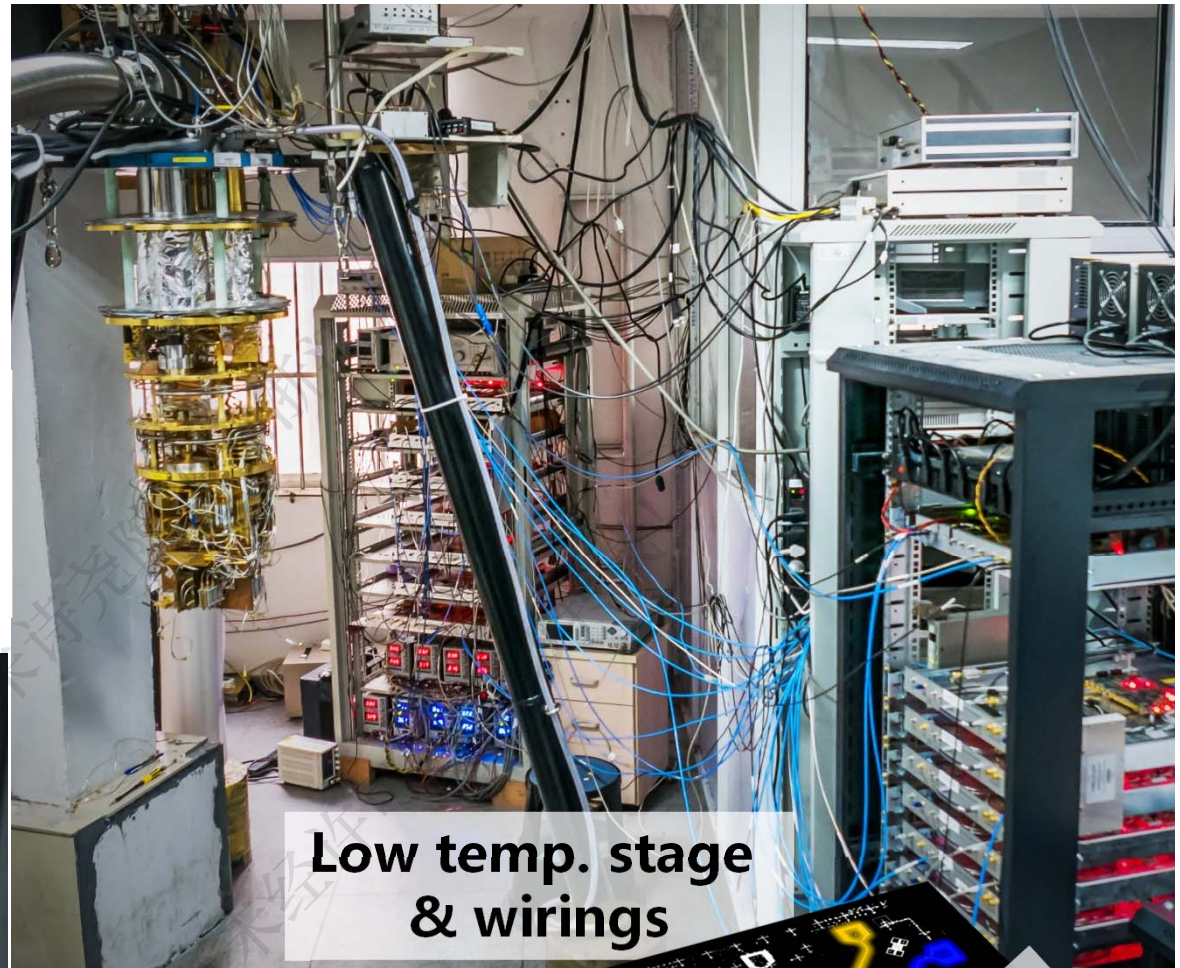
⇔ **Gate times**
10-100 ns

- ✓ Single-qubit gate fidelity above 99.9% [H. Wang group, PRL 119, 180511 (2017)].

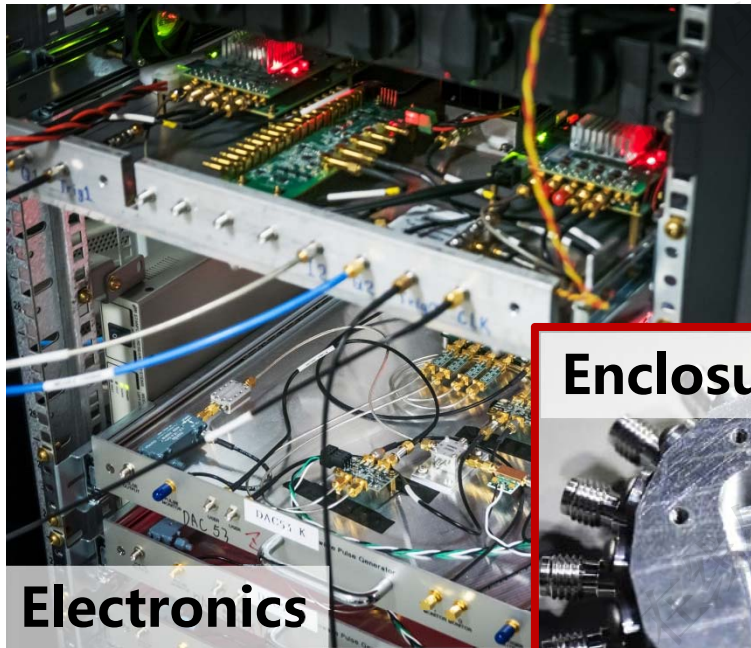
|ZJU-QLab>

- ✓ >20-Q device
- ✓ Multi-port enclosure
- ✓ Home-made electronics
- ✓ Large 10mK space

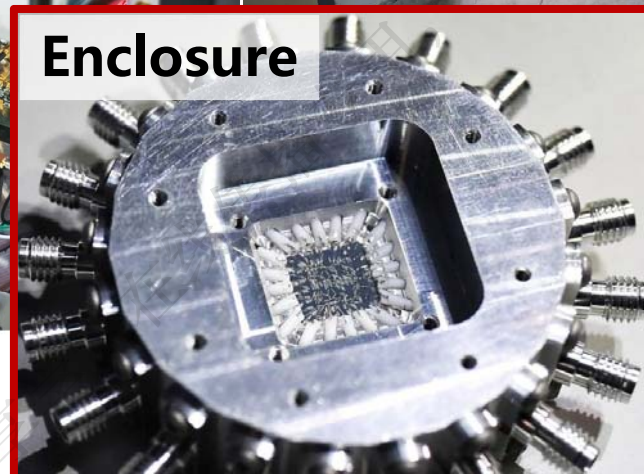
⇒ Powered by Q.E.



Low temp. stage
& wirings



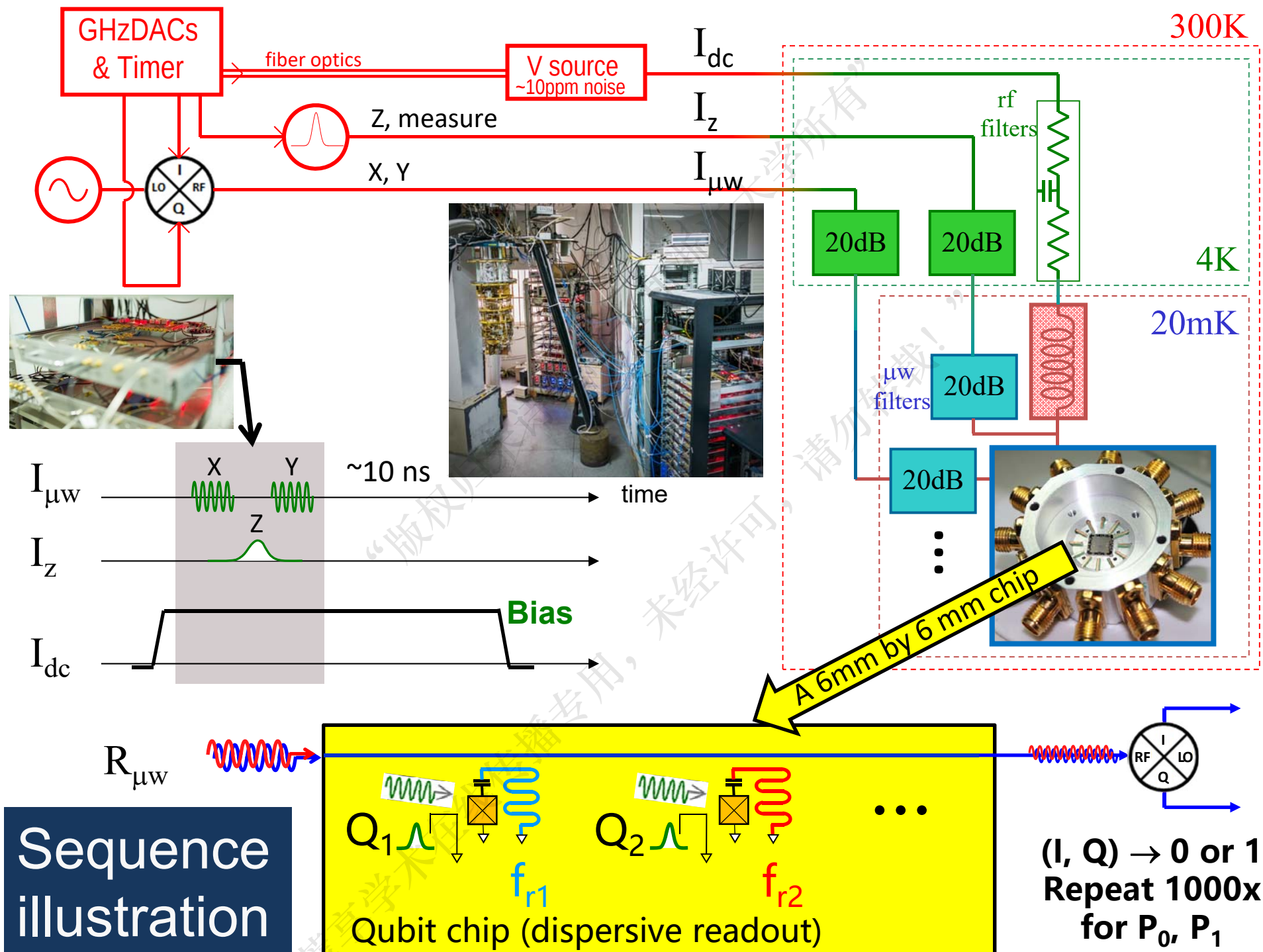
Electronics



Enclosure



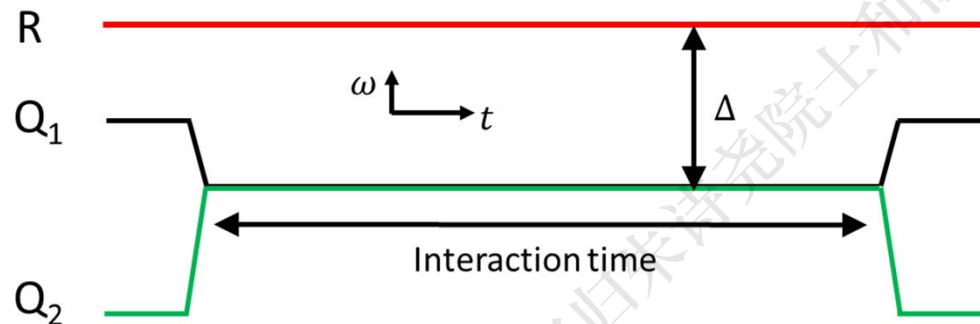
Device



Two-qubit interaction

- Interaction between two qubits enabled by bus resonator

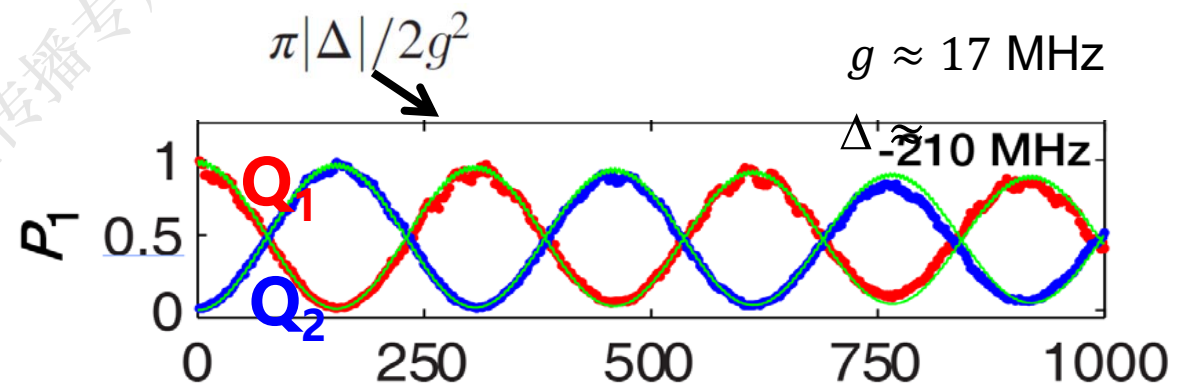
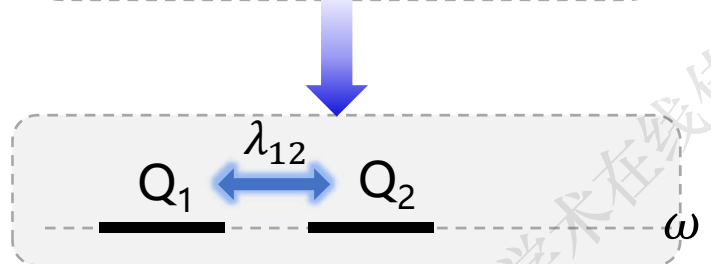
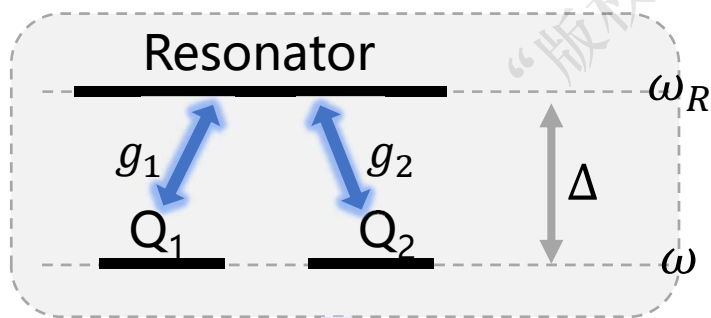
$$H = \sum_{j=1}^2 \omega_j \sigma_j^+ \sigma_j^- + \omega_R a^\dagger a + \sum_{j=1}^2 g_j (a \sigma_j^+ + a^\dagger \sigma_j^-),$$



- Resonator in vacuum state
- $\omega_1 = \omega_2 = \omega$, rotating frame
- $\Delta \gg g_1, g_2$

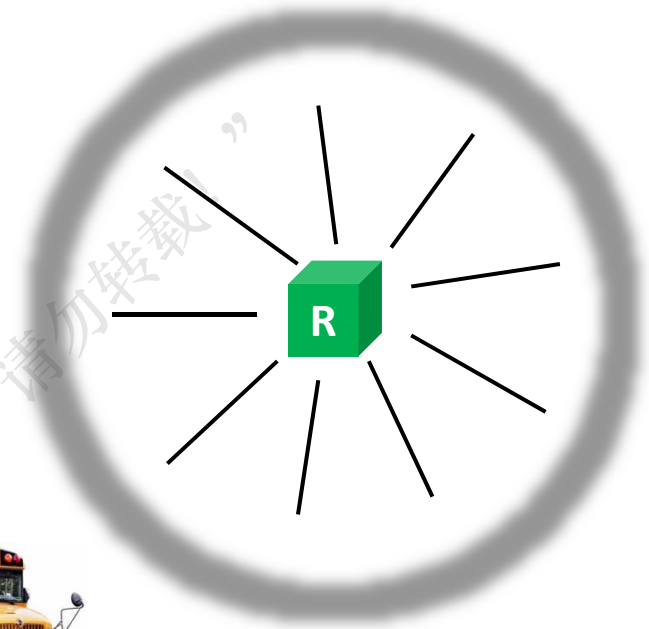
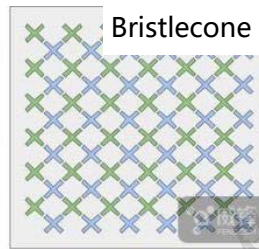
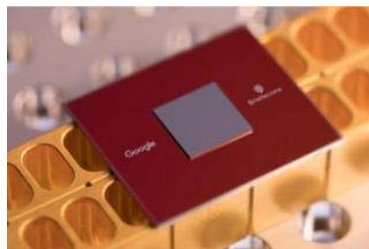
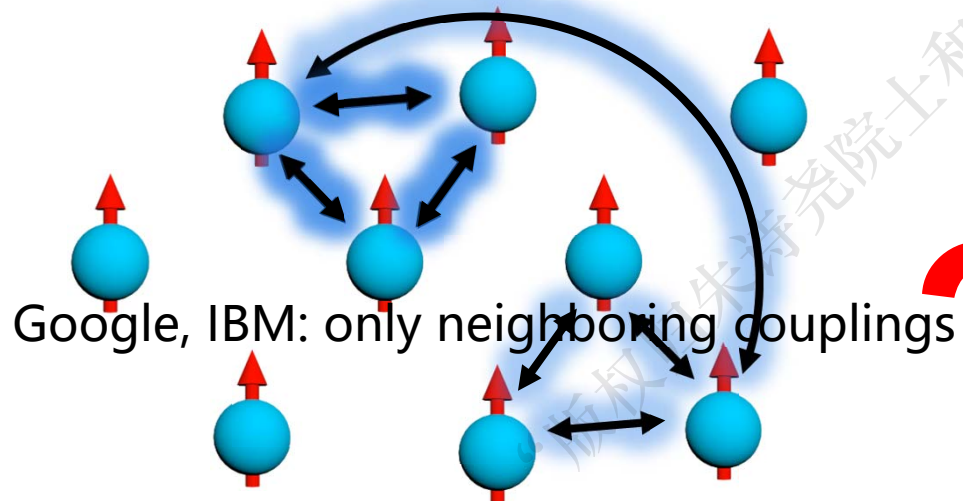
$$H_{eff} = \sum_{j=1}^2 \lambda_j \sigma_j^+ \sigma_j^- + \lambda_{12} (\sigma_1^+ \sigma_2^- + \sigma_1^- \sigma_2^+),$$

$$\text{where } \lambda_j = \frac{g_j^2}{\Delta}, \lambda_{12} = \frac{g_1 g_2}{\Delta}.$$



Circuit QED with many qubits (atoms)

- ✓ Wirings that can target various combinations of qubits

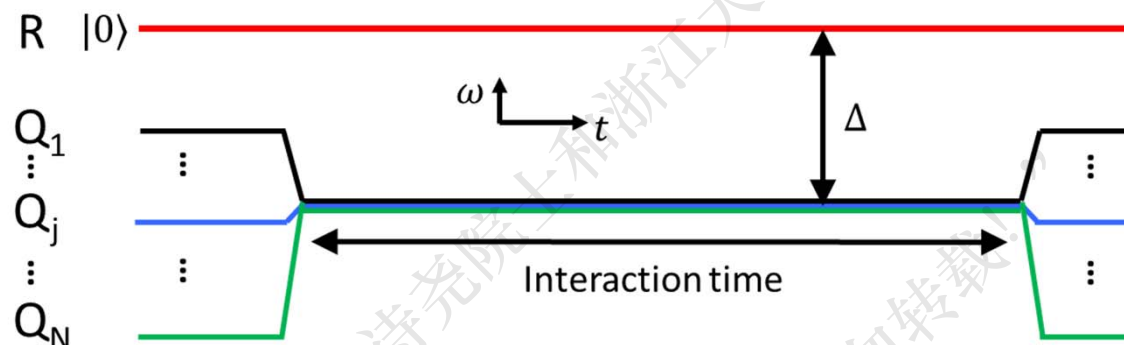


Interaction mediated
by bus resonator 

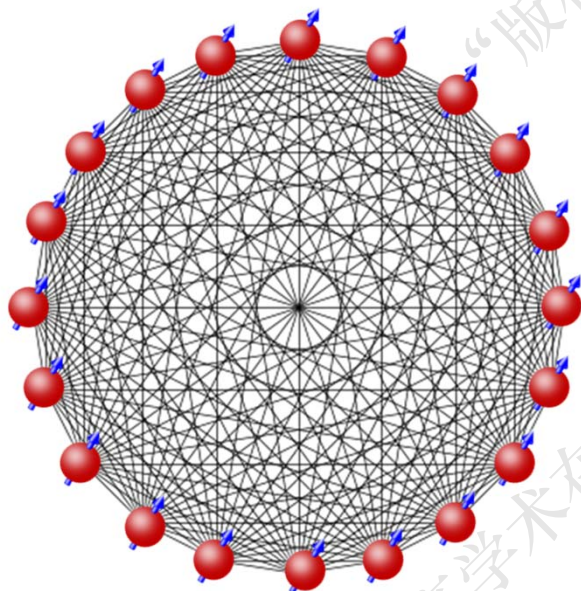
- ✓ Two-qubit and collective multiqubit interactions (*see next*)

Multiqubit interactions

- **Collective multiqubit interactions enabled by bus resonator**



- **Full-connectivity**

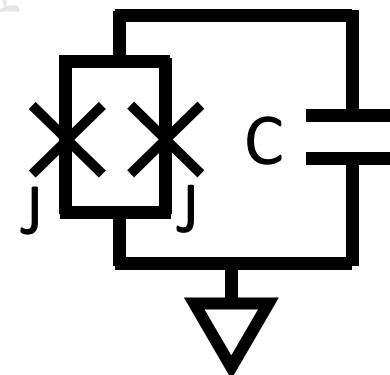
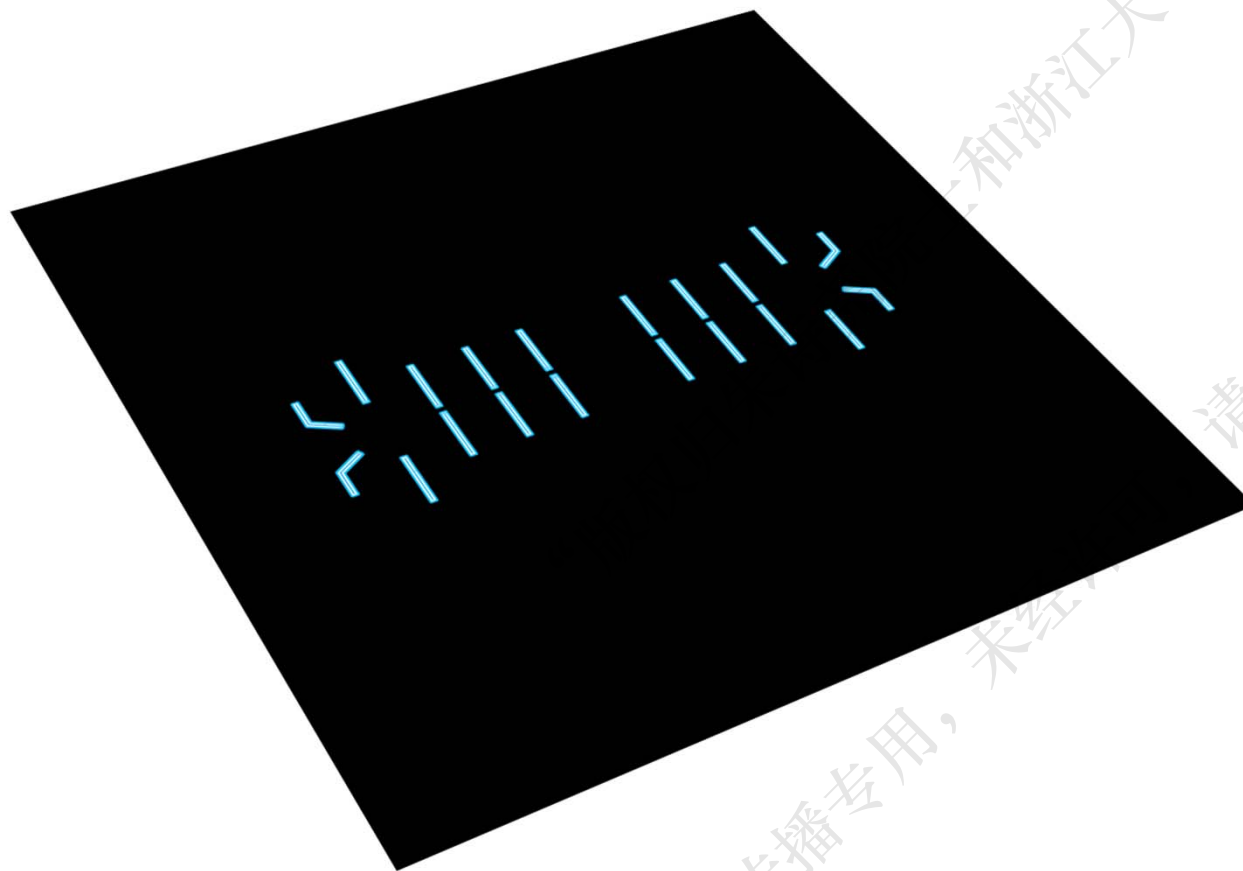


$$H_{eff} = \lambda \sum_{j=1}^N \sigma_j^+ \sigma_j^- + \sum_{i < j} \lambda (\sigma_i^+ \sigma_j^- + \sigma_j^+ \sigma_i^-)$$

- Uniform coupling strength g , $\lambda = g^2/\Delta$

Layout of the device

- Transmon Qubit $\times 20$



T_1 : $\sim 30 \mu\text{s}$

T_2 : $\sim 10 \mu\text{s}$ @ sweet point

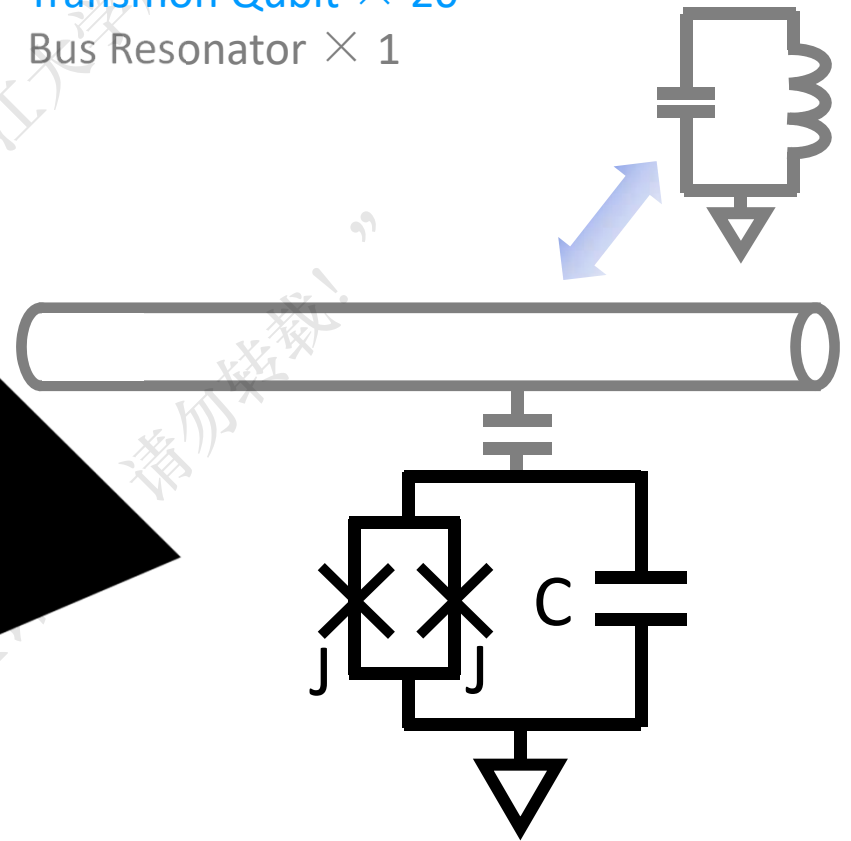
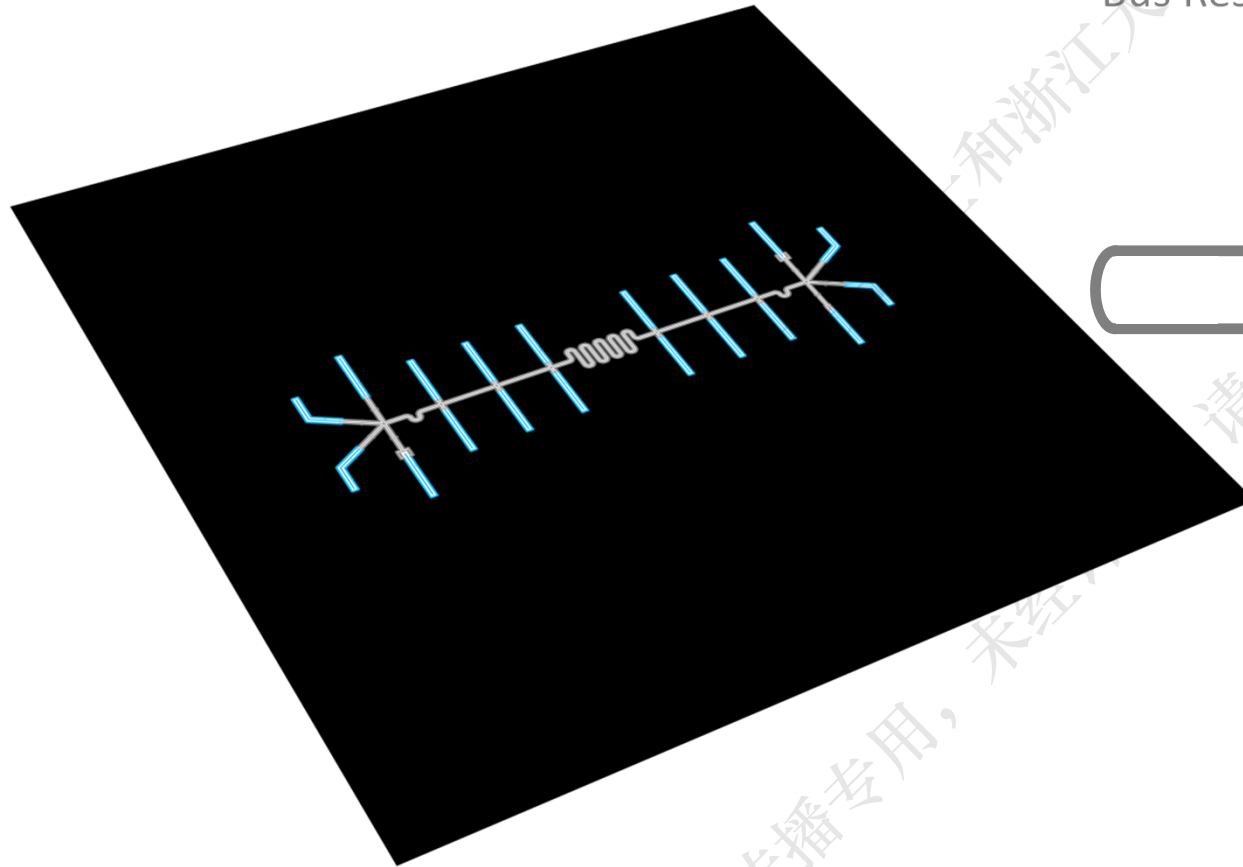
f_{sweet} : $\sim 5.7 \text{ GHz}$ (max. freq.)

$\eta/2\pi$: $\sim 240 \text{ MHz}$ (nonlin.)

E_J/E_C : ~ 90

Layout of the device

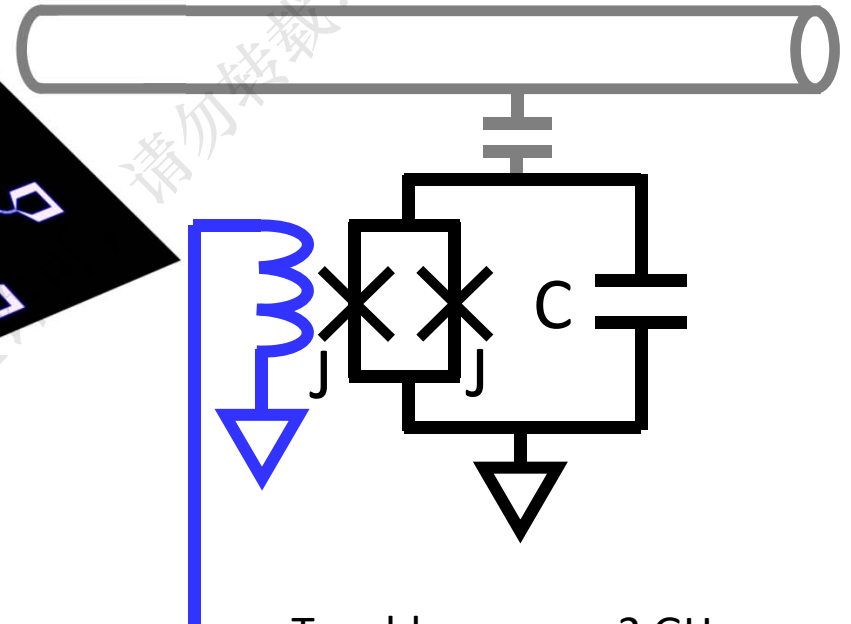
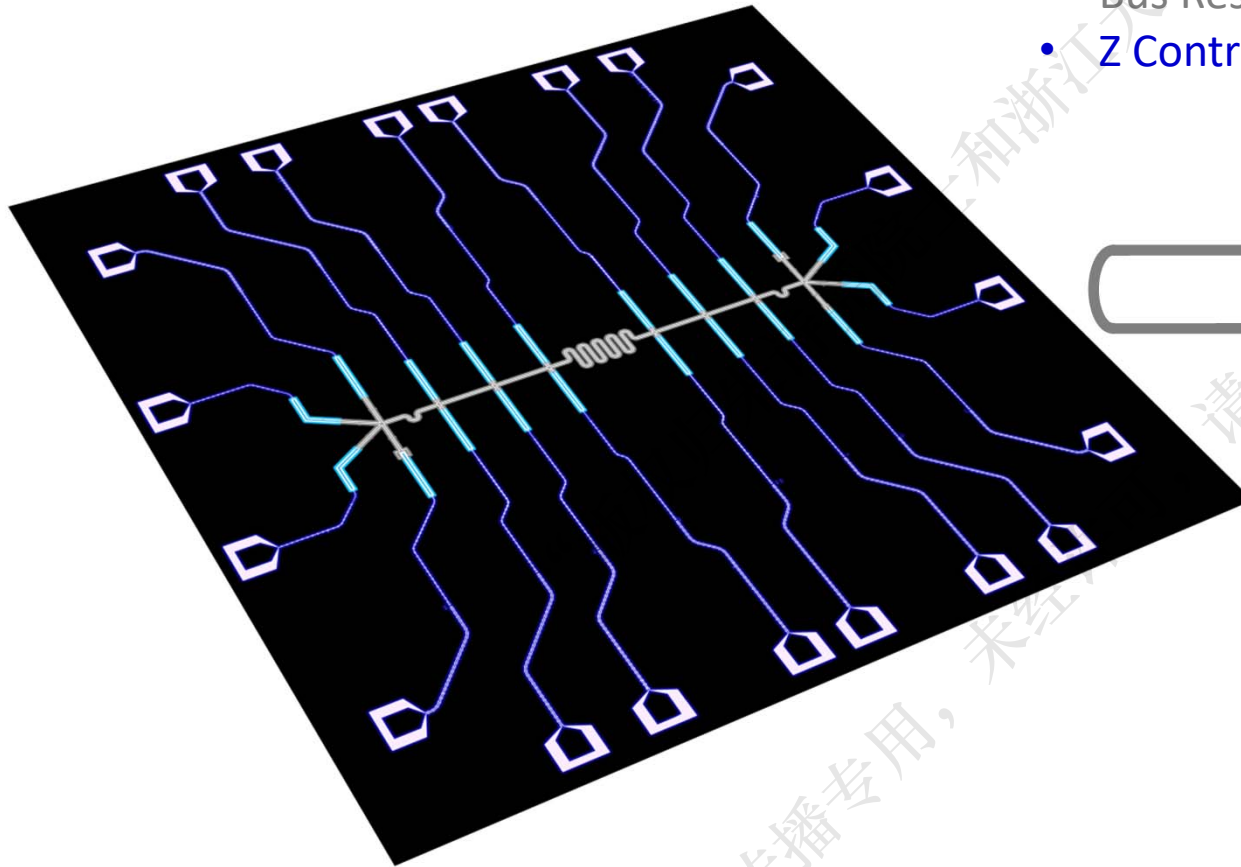
- Transmon Qubit $\times 20$
- Bus Resonator $\times 1$



$$f_R: \sim 5.5 \text{ GHz}$$
$$g/2\pi: \sim 27 \text{ MHz}$$

Layout of the device

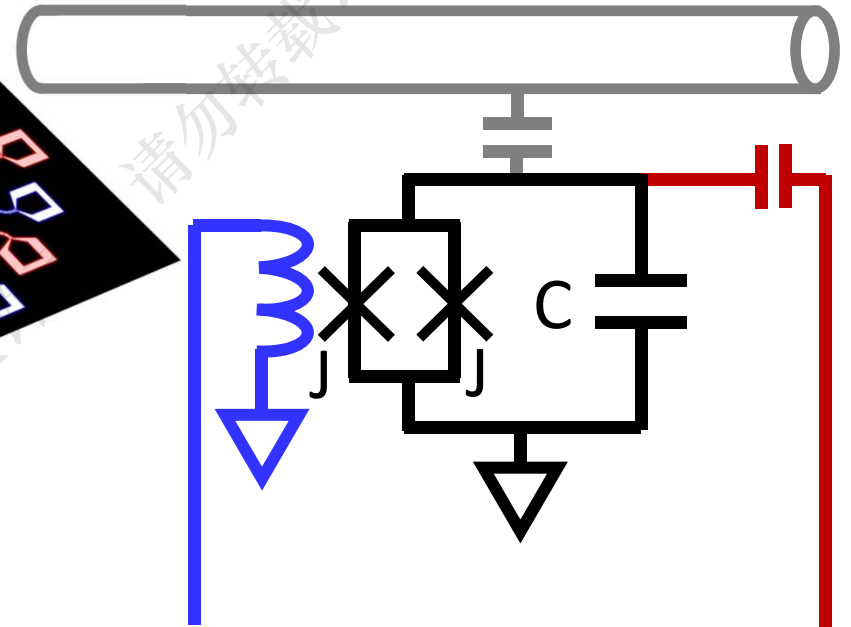
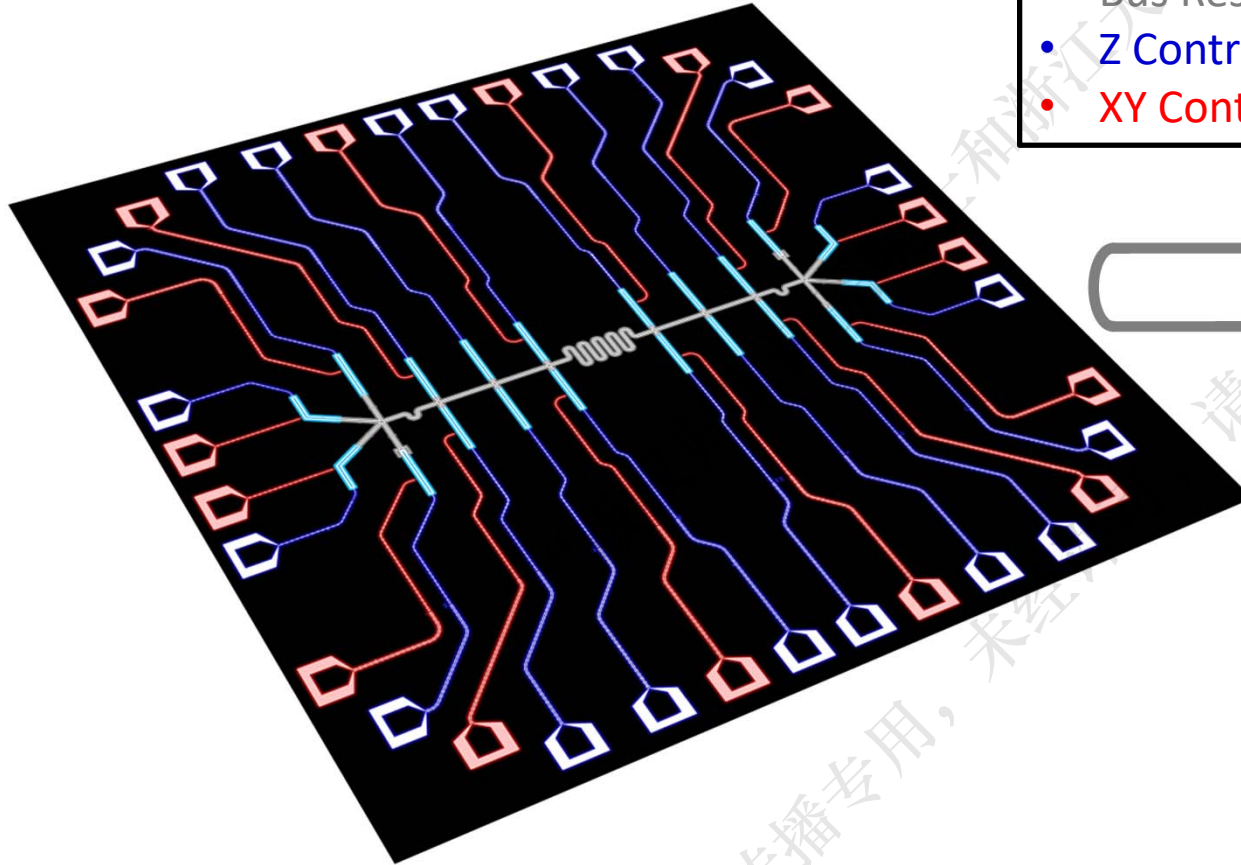
- Transmon Qubit $\times 20$
- Bus Resonator $\times 1$
- Z Control Line $\times 20$



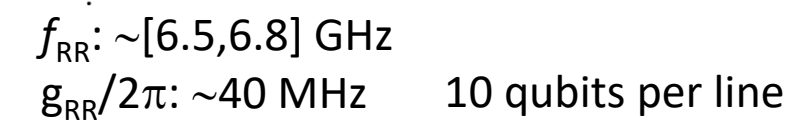
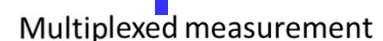
Tunable range: ~ 2 GHz

Layout of the device

- Transmon Qubit $\times 20$
- Bus Resonator $\times 1$
- Z Control Line $\times 20$
- XY Control Line $\times 16$

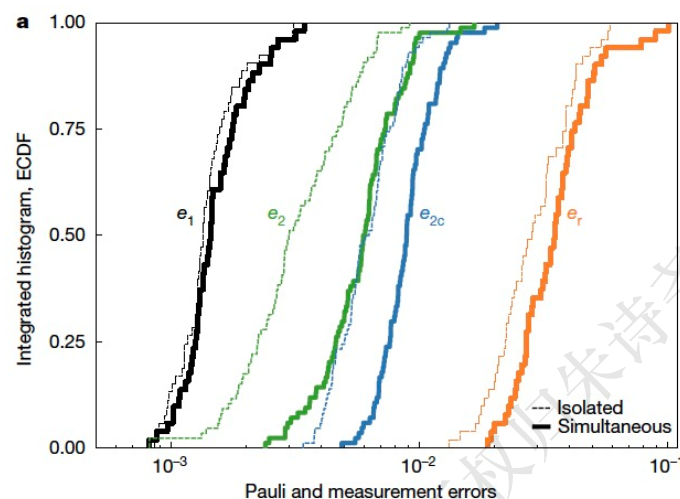


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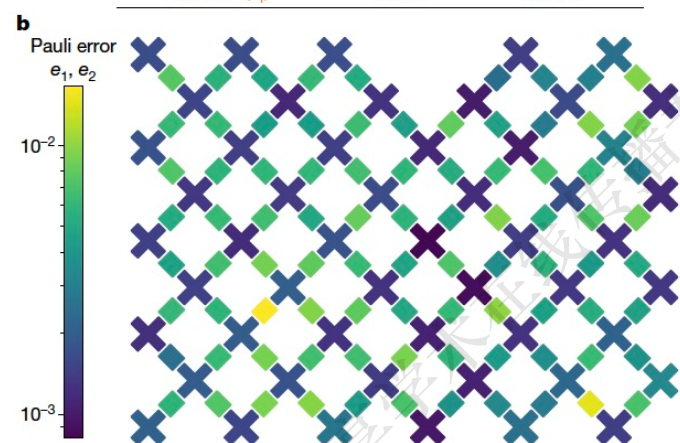


Current Quantum Computing Situation

Google 53-qubit sycamore quantum processor



Average error	Isolated	Simultaneous
Single-qubit (e_1)	0.15%	0.16%
Two-qubit (e_2)	0.36%	0.62%
Two-qubit, cycle (e_{2c})	0.65%	0.93%
Readout (e_r)	3.1%	3.8%



Quantum Supremacy

Matinis's group

Nature 574, 505 (2019)

Our Sycamore processor takes about **200 seconds** to sample one instance of a quantum circuit a million times—our benchmarks currently indicate that the equivalent task for a state-of-the-art classical supercomputer would take approximately **10,000 years**.

Fidelity $\sim 0.1\%$

(III) Quantum Optics Experiments with Superconducting-Qubits

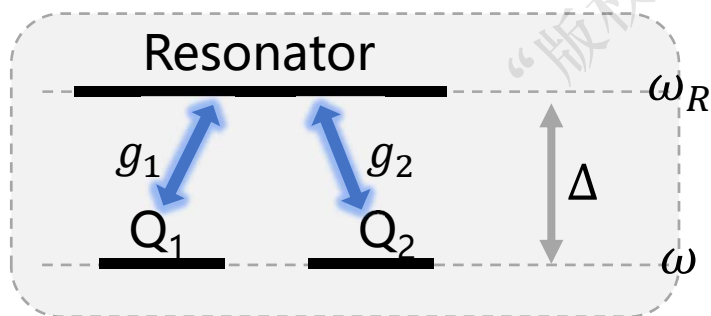
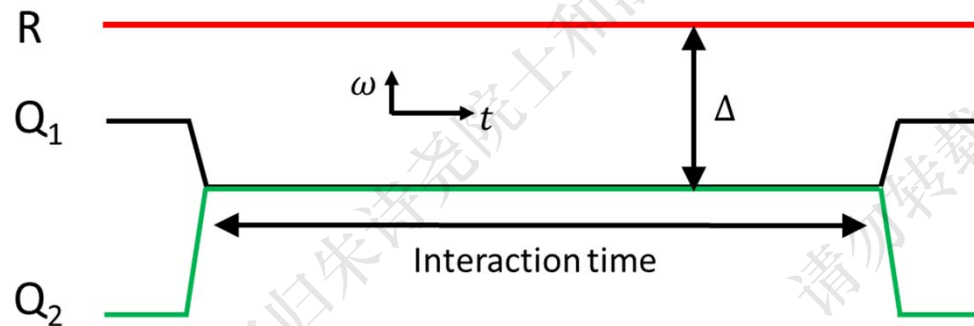


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“寇享学术在线传播专用，未经许可，请勿转载！”

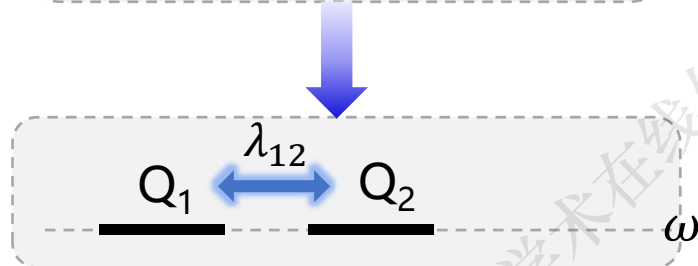
Two-qubit interaction

- Interaction between two qubits enabled by bus resonator

$$H = \sum_{j=1}^2 \omega_j \sigma_j^+ \sigma_j^- + \omega_R a^\dagger a + \sum_{j=1}^2 g_j (a \sigma_j^+ + a^\dagger \sigma_j^-),$$



- Resonator in vacuum state
- $\omega_1 = \omega_2 = \omega$, rotating frame
- $\Delta \gg g_1, g_2$



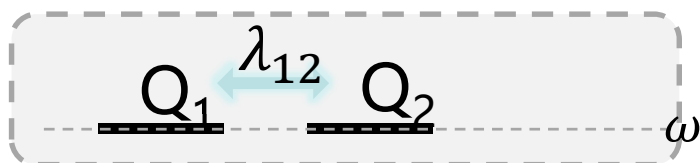
$$H_{eff} = \sum_{j=1}^2 \lambda_j \sigma_j^+ \sigma_j^- + \lambda_{12} (\sigma_1^+ \sigma_2^- + \sigma_1^- \sigma_2^+),$$

where $\lambda_j = \frac{g_j^2}{\Delta}$, $\lambda_{12} = \frac{g_1 g_2}{\Delta}$.

1. Quantum chiral rotation

- Anti-symmetric spin exchange interaction**

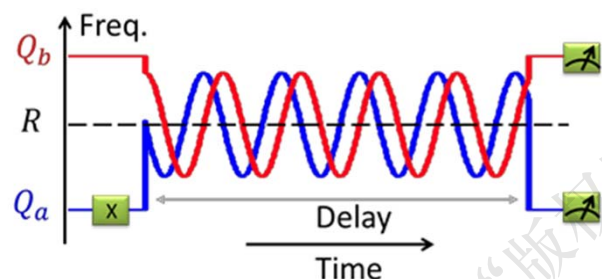
(Dzyaloshinskii-Moriya interaction)



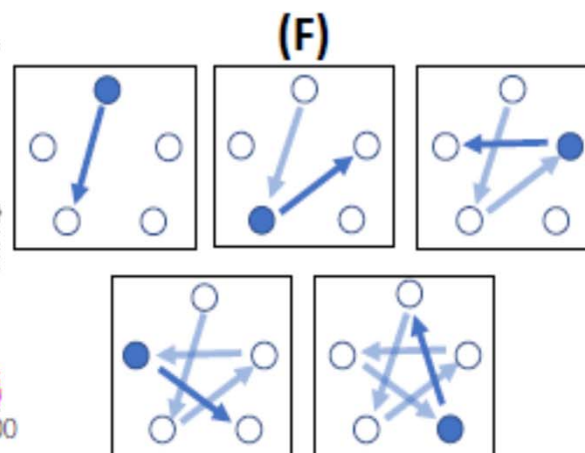
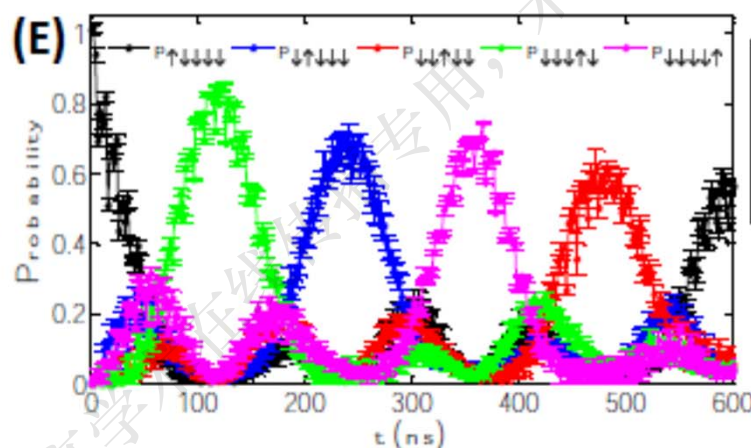
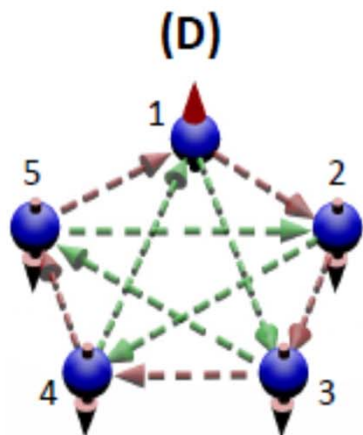
$$H_{eff} = \sum_{j=1}^2 \lambda_j \sigma_j^+ \sigma_j^- + i\lambda_{12}(\sigma_1^+ \sigma_2^- - \sigma_1^- \sigma_2^+),$$

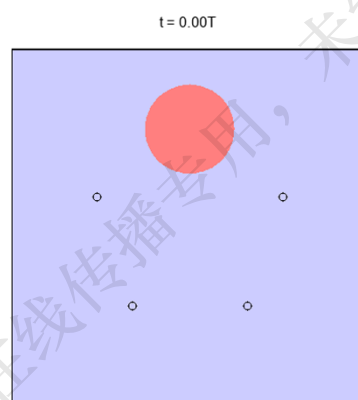
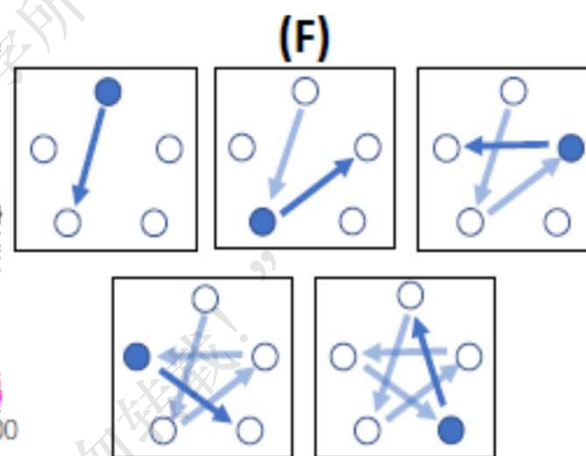
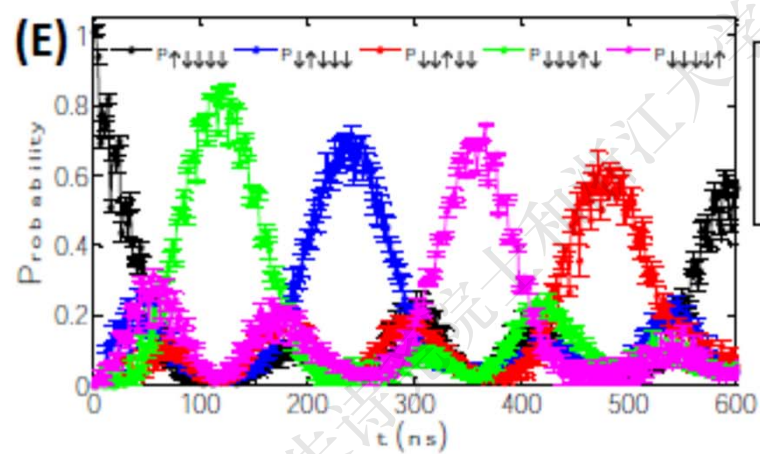
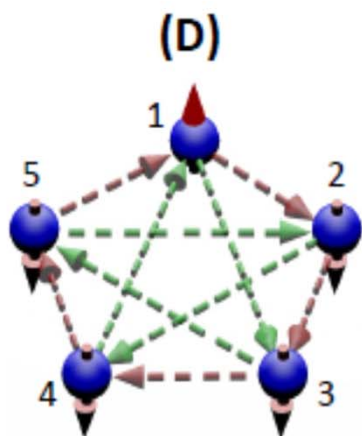
$$\lambda_{12} = \sum 2g_1 g_2 J_n^2(f) \sin n(\phi_2 - \phi_1) / nv.$$

$$\omega_j(t) = \omega_R + \Delta \cos(vt - \phi_j)$$



Wang et al, Nat. Phys. 15, 382 (2019)





2. 5-component Schrodinger cat states

Multiqubit interactions

$$H_{eff} = \lambda \sum_{j=1}^N \sigma_j^+ \sigma_j^- + \sum_{i < j} \lambda (\sigma_i^+ \sigma_j^- + \sigma_j^+ \sigma_i^-)$$

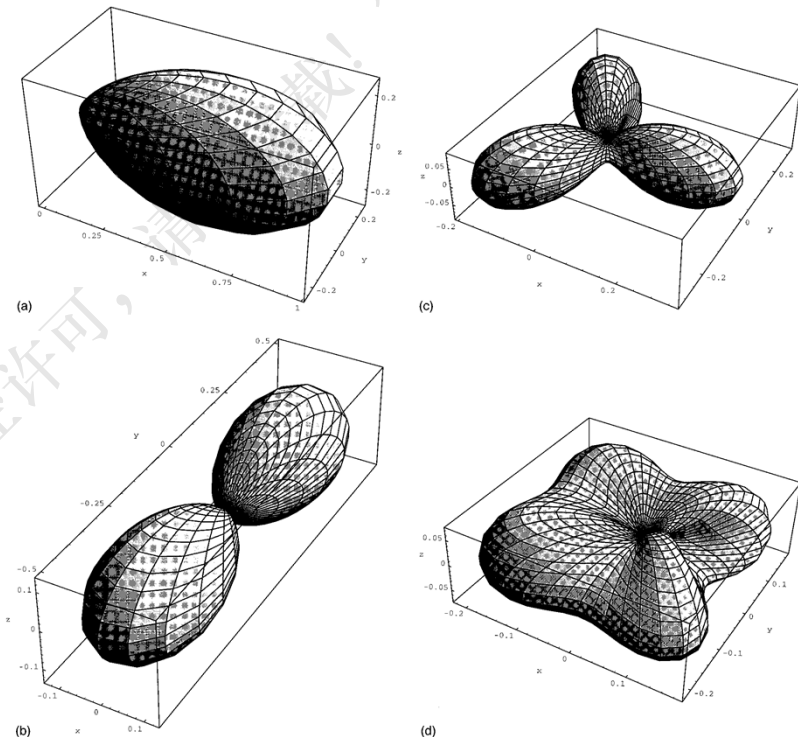
$$H_{eff} = \lambda S^+ S^-$$

$|s, m\rangle$ is the eigenstate of the system:
 $S^+ S^- |s, m\rangle = a |s, m\rangle$

$$\begin{aligned} S^+ S^- &= (S_x - iS_y)(S_x + iS_y) \\ &= S_x^2 + S_y^2 - S_z \\ &= S^2 - S_z^2 - S_z \end{aligned}$$

$$H_{eff}^S = \lambda [s(s+1) - S_z^2 - S_z]$$

One-axis twisting Hamiltonian in the subspace labeled by s



multi-component Schrodinger cat

Multi-qubit case [G. S. Agarwal, R. R. Puri, R. P. Singh, Phys. Rev. A 56, 2249 (1997)].

Schrödinger cat states

- Visualization and characterization**

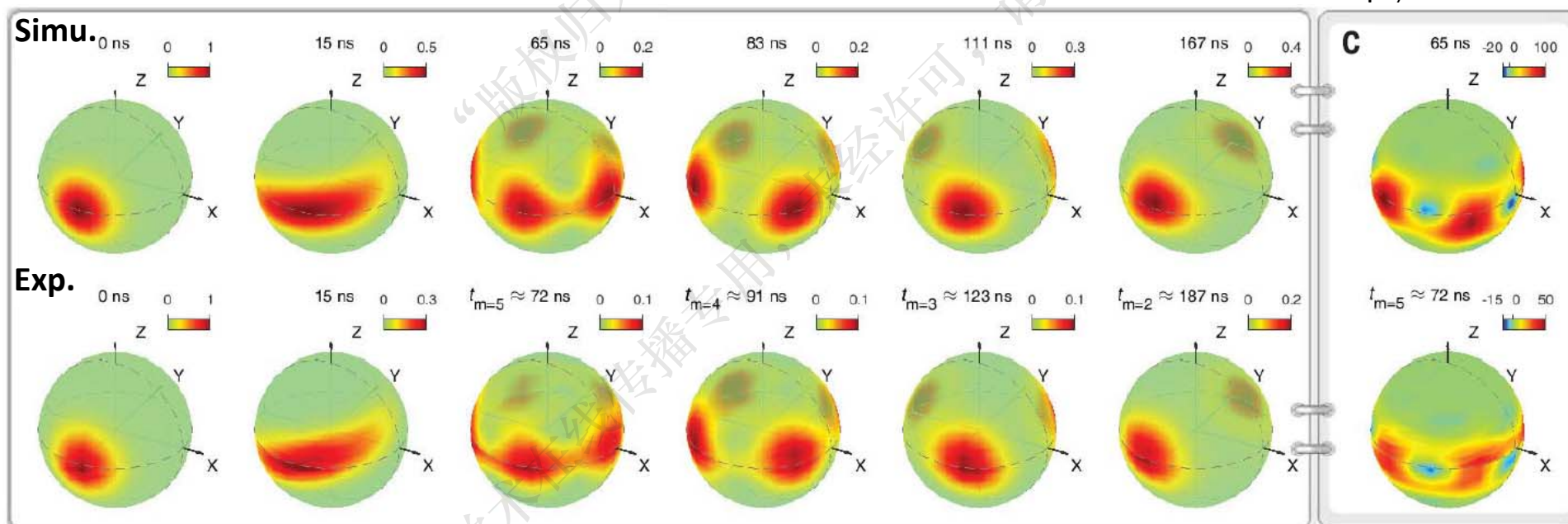
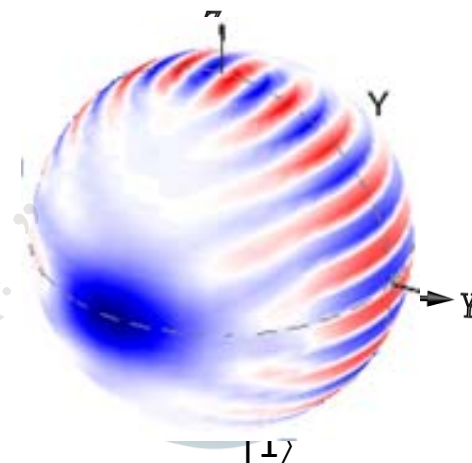
R. P. Rundle et al., Phys. Rev. A 96, 022117 (2017).

Sliced spin Wigner function:

$$W(\theta, \phi) = \text{Tr}(\hat{U}^\dagger(\theta, \phi) \rho \hat{U}(\theta, \phi) \hat{\Pi})$$

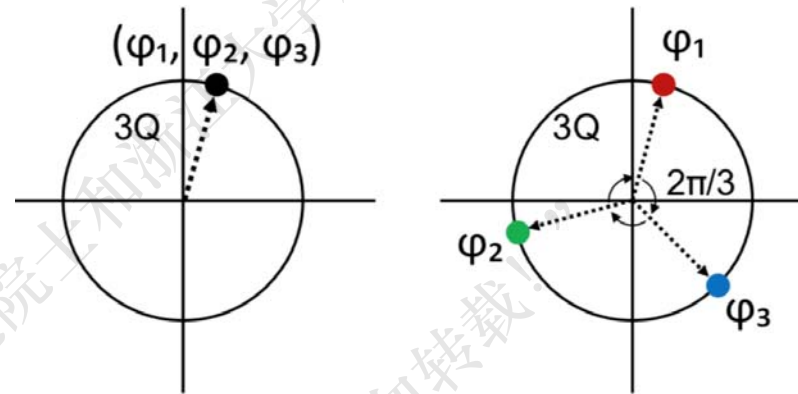
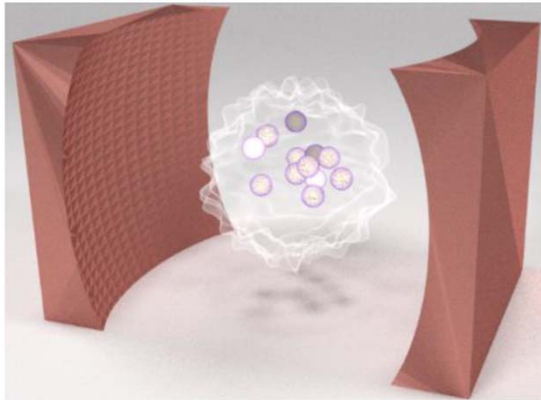
$$\hat{U}(\theta, \phi) = \bigotimes_{j=1}^N \hat{u}_j(\theta, \phi)$$

$$\hat{\Pi} = \bigotimes_{j=1}^N (1 - \sqrt{3}\sigma_{z,j}) \longrightarrow \text{extended parity operator}$$

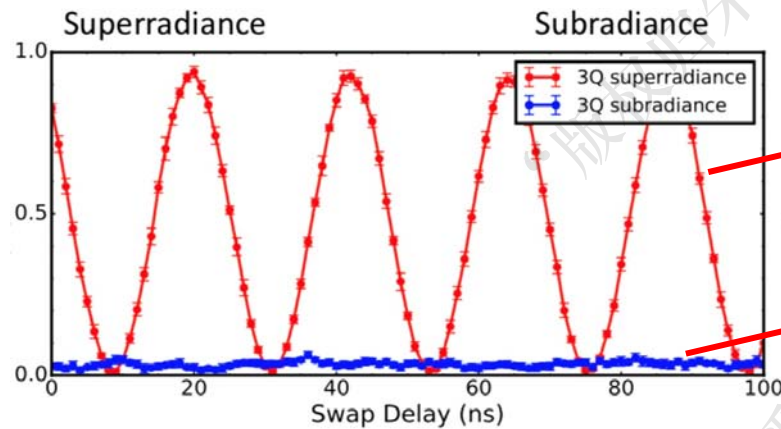


Song et al, Science 365, 574-577 (2019).

3. Control of Superradiant and subradiant states



superradiant and subradiant states



No. of cavity photons

$$|\psi_{super}\rangle = \frac{1}{\sqrt{3}} (|100\rangle + |010\rangle + |001\rangle)$$

$$|\psi_{sub}\rangle = \frac{1}{\sqrt{3}} (|001\rangle + e^{i2\pi/3}|010\rangle + e^{i4\pi/3}|001\rangle)$$

$$\langle 000 | \sum \sigma_j^- | \psi_{super} \rangle = \sqrt{3}$$

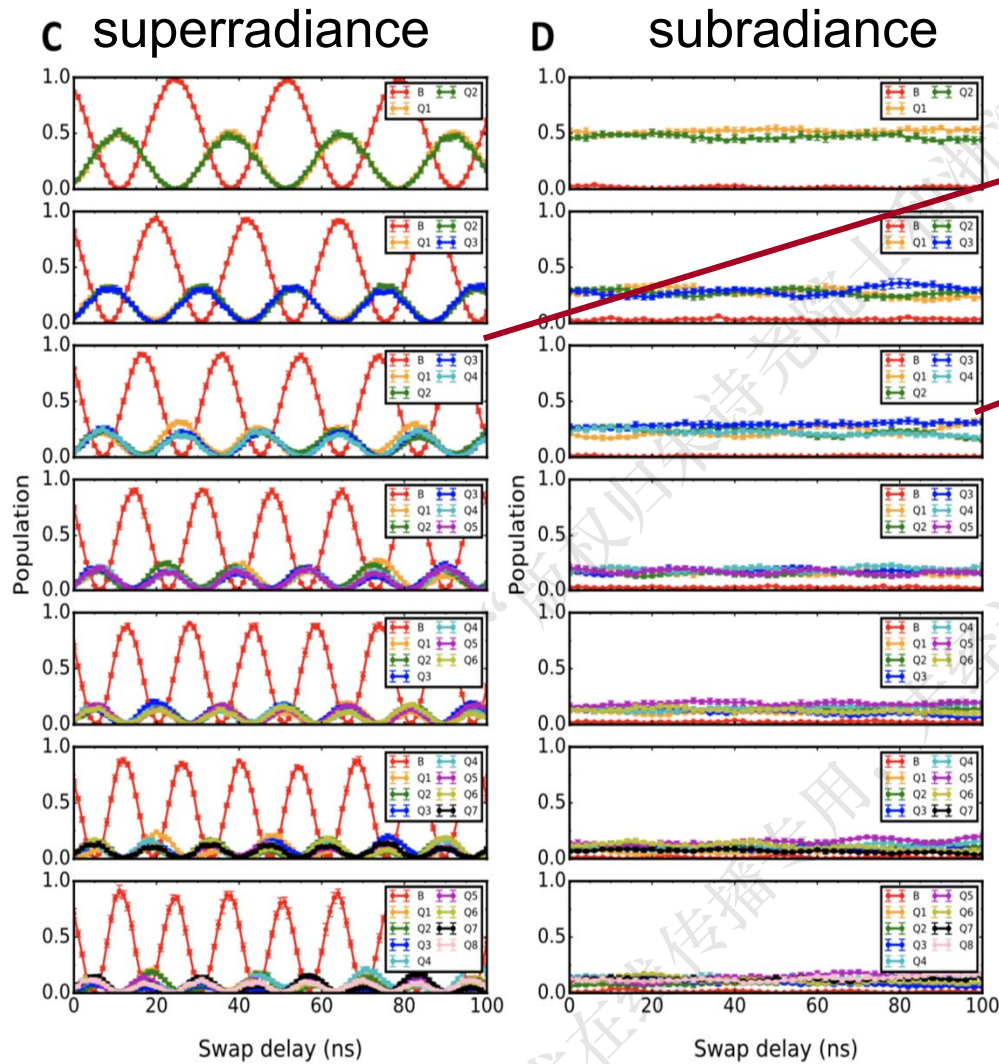
$$\langle 000 | \sum \sigma_j^- | \psi_{sub} \rangle = 0$$

Scully, PRL 115, 243602 (2015)

Vetter et al, Phys. Scr. 91, 023007 (2016)

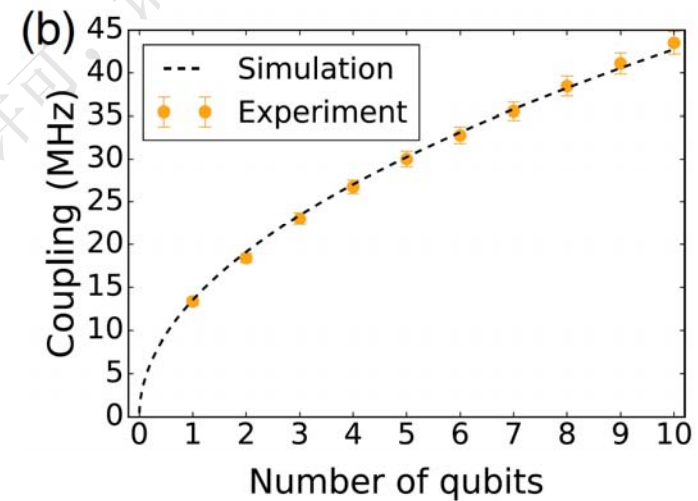
Wang et al, PRL 124, 013601 (2020)

Superradiant and subradiant state



$$|\psi_{super}\rangle = \frac{1}{\sqrt{N}} \sum_j \sigma_j^+ |g_1, \dots, g_N\rangle$$

$$|\psi_{sub}\rangle = \frac{1}{\sqrt{N}} \sum_j e^{i2\pi j/N} \sigma_j^+ |g_1, \dots, g_N\rangle$$



$\sqrt{N}$ -scaling

Wang et al, PRL 124, 013601 (2020)

3. Two-atom excited by entangled photons

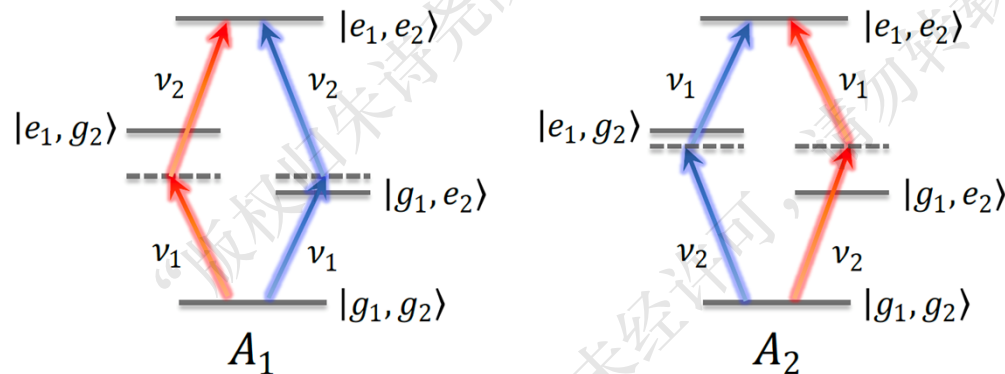
VOLUME 93, NUMBER 9

PHYSICAL REVIEW LETTERS

week ending
27 AUGUST 2004

Inducing Disallowed Two-Atom Transitions with Temporally Entangled Photons

Ashok Muthukrishnan,¹ Girish S. Agarwal,² and Marlan O. Scully^{1,3}

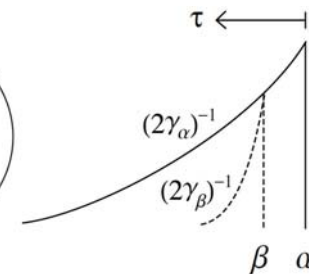
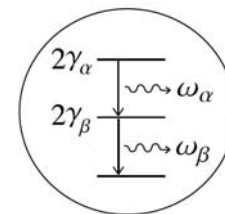


$$A_1 = \frac{\kappa_{22}\kappa_{11}}{\nu_1 - \omega_1} + \frac{\kappa_{12}\kappa_{21}}{\nu_1 - \omega_2},$$

$$A_2 = \frac{\kappa_{21}\kappa_{12}}{\nu_2 - \omega_1} + \frac{\kappa_{11}\kappa_{22}}{\nu_2 - \omega_2},$$

$$|A_1 + A_2|^2 = 0$$

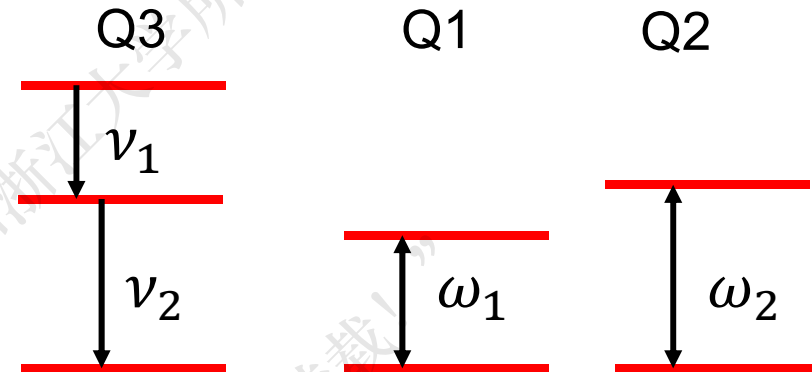
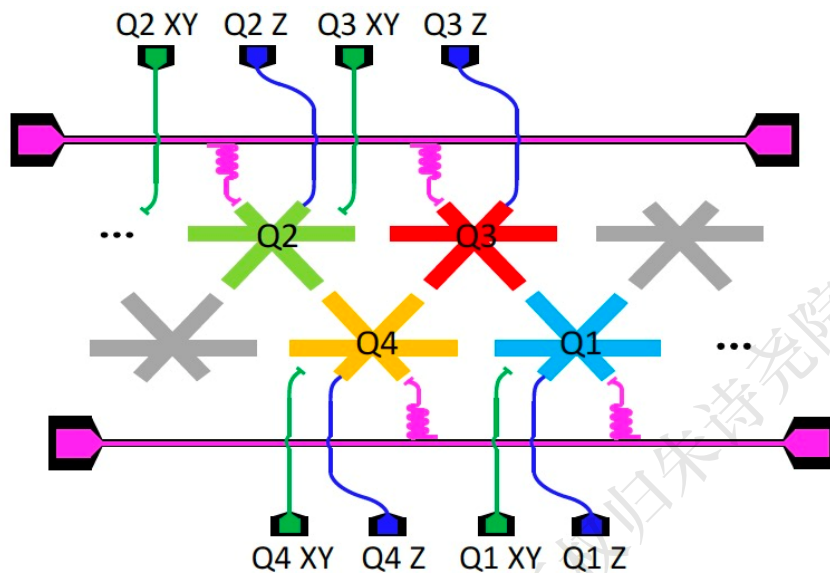
Independent photons



$$|A_1|^2 \neq 0$$

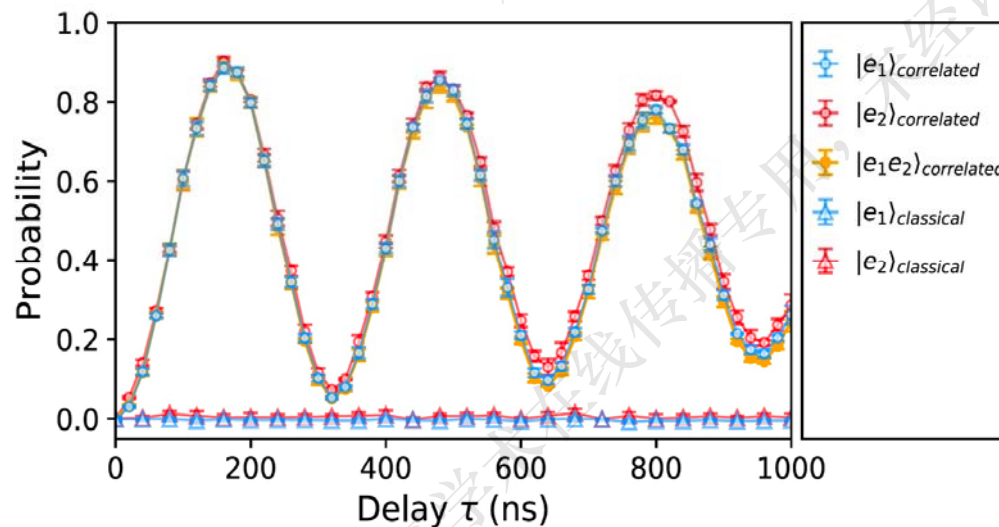
entangled photons

Two-atom excited by entangled photons



$$\nu_1 + \nu_2 = \omega_1 + \omega_2,$$

$$\nu_i \neq \omega_j \text{ for all } i, j = 1, 2$$



Effective 3-body int. H.

$$H_{\text{eff}} = A_1 \sigma_1^+ \sigma_2^+ \Xi_3^- + h.c.$$

Ren et al, RPL 125, 133601 (2020)

Outlook



1. Devices and new physics

Experiments with more than 20 qubits is going on

- New quantum device: quantum circulator etc.
- Topological physics of quantized light

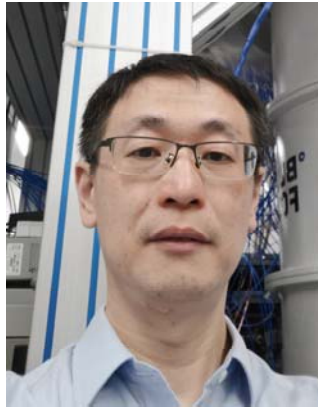
2. Quantum simulation

- Many-body interaction: $\sigma_1^- \sigma_2^+ \sigma_3^+ \sigma_4^+ \sigma_5^+$ etc
- Quantum origin of chiral molecules
- Simulation of energy transfer in photo-synthetic processes

3. Applications in quantum chemistry

- Ground state energy estimation of molecules

Collaborators



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Wenhui Ren, Hekang Li, Zhen Wang

IOP: Dongning Zheng, Xiaohui Song, Heng Fan

USTC: Xiaobo Zhu, Hui Deng

Thanks

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